

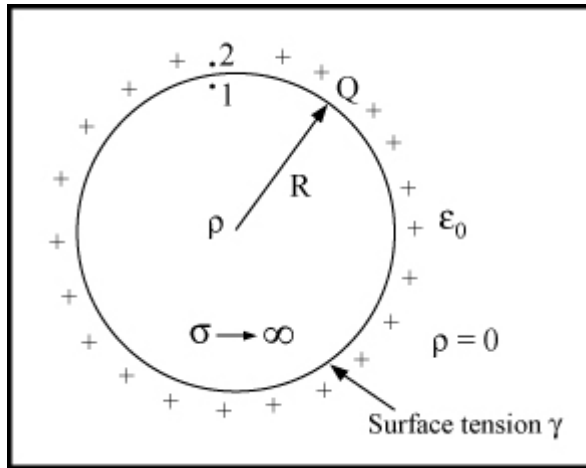
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6.642 Continuum Electromechanics
Fall 2008

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Lecture 10: Stability of a Perfectly Conducting Spherical Drop (Rayleigh's Limit)
 Continuum Electromechanics (Melcher) – Section 8.13

I. Geometry-Equilibrium



$$E_r = \frac{Q}{4\pi\epsilon_0 r^2} = E_0 \left(\frac{R}{r}\right)^2 \text{ for } r < R$$

$$E_0 = \frac{Q}{4\pi\epsilon_0 R^2}$$

$$\bar{T}_s = \frac{-2\gamma}{R} \bar{n}$$

$$P_2 = 0 \text{ (air or vacuum, density=0)}$$

$$\bar{E}_1 = 0 \text{ (perfectly conducting drop)}$$

$$P_1 - \frac{2\gamma}{R} + \frac{1}{2} \epsilon_0 E_r (r=R)^2 = 0$$

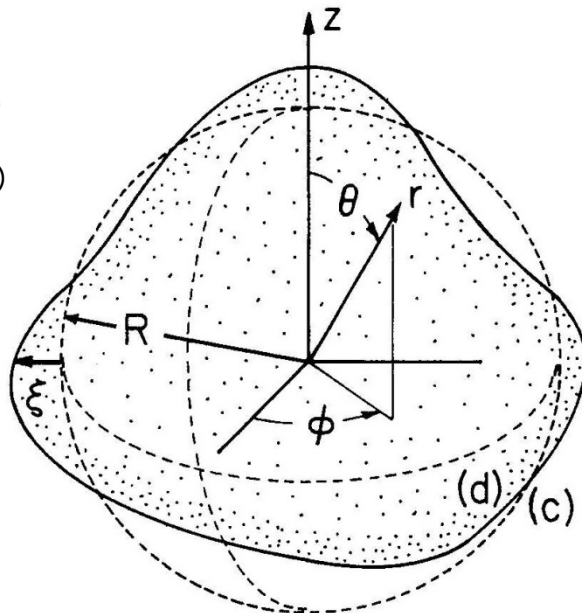
II. Perturbations

$$\xi = \text{Re} \left[\hat{\xi} P_n^m (\cos \theta) e^{j(\omega t - m\phi)} \right]$$

$$F = \xi(\theta, \phi, t) - r + R$$

$$\bar{n} = \frac{\nabla F}{|\nabla F|} = \bar{i}_r - \frac{1}{R} \frac{\partial \xi}{\partial \theta} \bar{i}_\theta - \frac{1}{R \sin \theta} \frac{\partial \xi}{\partial \phi} \bar{i}_\phi$$

$$\bar{T}_s = -\gamma \left(\frac{1}{R_1} + \frac{1}{R_2} \right) \bar{n} = -\gamma (\nabla \cdot \bar{n}) \bar{n}$$



Spherically symmetric equilibrium for a drop having total charge q uniformly distributed over its surface.

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$$\hat{T}_{sr} = -\frac{\gamma}{R^2}(n-1)(n+2)\hat{\xi}$$

$$\hat{e}_{r1} = 0 \text{ (perfectly conducting drop)}$$

$$\hat{e}_{r2} = \frac{(n+1)\hat{\Phi}_2}{R}$$

$$\hat{P}_2 = 0 \text{ (air or vacuum, density=0)}$$

$$\hat{P}_1 = \frac{\rho\omega^2 R}{n}\hat{\xi}$$

III. Electric Field Boundary Conditions

$$\bar{n} \times \bar{E}|_{r=R+\xi} = 0 \Rightarrow \left[\bar{i}_r - \frac{1}{R} \frac{\partial \xi}{\partial \theta} \bar{i}_\theta - \frac{1}{R \sin \theta} \frac{\partial \xi}{\partial \phi} \bar{i}_\phi \right] \times \left[(E_0(R+\xi) + e_r(R)) \bar{i}_r + e_\theta(R) \bar{i}_\theta + e_\phi(R) \bar{i}_\phi \right] = 0$$

$$\bar{i}_\phi \left[e_\theta + \frac{E_0}{R}(R) \frac{\partial \xi}{\partial \theta} \right] + \bar{i}_\theta \left[-e_\phi - \frac{1}{R \sin \theta} E_0(R) \frac{\partial \xi}{\partial \phi} \right] = 0$$

$$\left. \begin{aligned} e_\theta &= -\frac{1}{R} \frac{\partial \Phi_2}{\partial \theta} = -\frac{E_0(R)}{R} \frac{\partial \xi}{\partial \theta} \\ e_\phi &= -\frac{1}{R \sin \theta} \frac{\partial \Phi_2}{\partial \phi} = -\frac{E_0(R)}{R \sin \theta} \frac{\partial \xi}{\partial \phi} \end{aligned} \right\} \Rightarrow \hat{\Phi}_2 = E_0(R) \hat{\xi}$$

IV. Interfacial Force Balance Boundary Conditions

$$P_{10} - P_{20} + \hat{P}'_1 - \hat{P}'_2 + T_{rj}n_j - \frac{2\gamma}{R} - \frac{\gamma}{R^2}(n-1)(n+2)\hat{\xi} = 0$$

$$T_{ij} = \varepsilon_0 E_i E_j - \frac{1}{2} \delta_{ij} \varepsilon_0 E_k E_k$$

$$T_{rj}n_j = T_{rr}n_r + T_{r\theta}n_\theta + T_{r\phi}n_\phi$$

$$\begin{aligned} T_{rr} &= \frac{1}{2} \varepsilon_0 \left[E_r^2 - \underbrace{e_\theta^2 + e_\phi^2}_{\text{Second order}} \right] = \frac{1}{2} \varepsilon_0 \left[\left(E_r(r=R) + \frac{dE_r}{dr} \Big|_{r=R} \xi + e_r \right)^2 \right] \\ &= \frac{1}{2} \varepsilon_0 \left[E_0^2 + 2 \left(\frac{dE_r}{dr} \Big|_{r=R} \xi + e_r \right) E_0 \right] \end{aligned}$$

$$\left. \begin{aligned} T_{r0} &= \varepsilon_0 E_r e_0 \Rightarrow T_{r0} n_0 \\ T_{r\phi} &= \varepsilon_0 E_r e_\phi \Rightarrow T_{r\phi} n_\phi \end{aligned} \right\} \text{second order}$$

$$T_{rj} n_j = \frac{1}{2} \varepsilon_0 E_0^2 + \varepsilon_0 \left[\frac{dE_r}{dr} \Big|_{r=R} \xi + e_{r2} \right] E_0$$

$$E_r = \frac{E_0 R^2}{r^2} \Rightarrow \frac{dE_r}{dr} \Big|_{r=R} = \frac{-2E_0 R^2}{R^3} = \frac{-2E_0}{R}$$

$$\underbrace{P_{10} - P_{20} + \frac{1}{2} \varepsilon_0 E_0^2 - \frac{2\gamma}{R}}_0 + \hat{P}_1 - \hat{P}_2 + \varepsilon_0 E_0 \left[\frac{-2E_0}{R} \hat{\xi} + \hat{e}_{r2} \right] - \frac{\gamma}{R^2} (n-1)(n+2) \hat{\xi} = 0$$

=0 (Equilibrium)

V. Dispersion Relation

$$\frac{\rho \omega^2 R}{n} \hat{\xi} + \varepsilon_0 \left[\frac{-2E_0^2}{R} + \frac{(n+1)E_0^2}{R} \right] \hat{\xi} - \frac{\gamma}{R^2} (n-1)(n+2) \hat{\xi} = 0$$

$$\rho \omega^2 R = \frac{n(n-1)}{R} \left[\frac{\gamma}{R} (n+2) - \varepsilon_0 E_0^2 \right]$$

$$\text{First unstable when } n=2 \Rightarrow \varepsilon_0 E_0^2 = \frac{4\gamma}{R} = \varepsilon_0 \left(\frac{Q}{4\pi \varepsilon_0 R^2} \right)^2$$

$$Q = 8\pi \sqrt{\varepsilon_0 \gamma R^3} \quad [\text{Rayleigh Limit}]$$