

Plate Bending

not so long plate

previously have shown: $M := -D \cdot \frac{d^2}{dx^2} w$. this was for single axis bending.

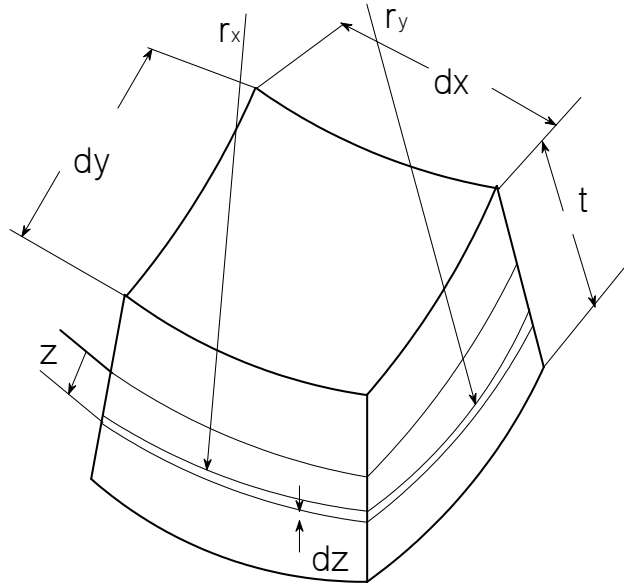
this relationship holds for the partial derivative in the respective direction fo both x and y;

assumptions:

plane cross section remains plane

small deflections $w_{\max} < 3/4 t$

stress < yield



$$w = w(x, y)$$

$$\frac{1}{R_x} = -\frac{d^2}{dx^2} w(x, y) \quad \frac{1}{R_y} = -\frac{d^2}{dy^2} w(x, y)$$

we are making no statement with respect to ε_y (or any other ε) as we did in long plate.

$$\varepsilon_x := \frac{1}{R_x} \cdot z \quad \varepsilon_x := -z \cdot \frac{d^2}{dx^2} w(x, y) \quad \varepsilon_y := \frac{1}{R_y} \cdot z \quad \varepsilon_y := -z \cdot \frac{d^2}{dy^2} w(x, y)$$

$$\varepsilon_x := \frac{\sigma_x}{E} - \frac{\nu \cdot \sigma_y}{E} \quad \varepsilon_y := \frac{\sigma_y}{E} - \frac{\nu \cdot \sigma_x}{E} \quad \nu \cdot \varepsilon_y \rightarrow$$

$$\varepsilon_x + \nu \cdot \varepsilon_y \text{ collect, } \sigma_x, E \rightarrow \quad \Rightarrow \quad \sigma_x := \frac{E}{1 - \nu^2} \cdot (\varepsilon_x + \nu \cdot \varepsilon_y)$$

substituting

into

$$\varepsilon_y := -z \frac{d^2}{dy^2} w(x, y) \quad \varepsilon_x := -z \frac{d^2}{dx^2} w(x, y) \quad \sigma_x := \frac{E}{1 - \nu^2} (\varepsilon_x + \nu \varepsilon_y)$$

$$\sigma_x \rightarrow \frac{E}{1 - \nu^2} \left(-z \frac{d}{dx} \frac{d}{dx} w(x, y) - \nu \cdot z \frac{d}{dy} \frac{d}{dy} w(x, y) \right) \text{ or ... } \sigma_x(z) := -z \frac{E}{1 - \nu^2} \left(\frac{d^2}{dx^2} w(x, y) + \nu \frac{d^2}{dy^2} w(x, y) \right)$$

$$\text{similarly: } \sigma_y := -z \frac{E}{1 - \nu^2} \left(\frac{d^2}{dy^2} w(x, y) + \nu \frac{d^2}{dx^2} w(x, y) \right)$$

as in bending (applied in each of x and y direction): see figure Hughes 9.3 (below)

the lower case m denotes moment per unit length

note that the designation is changed: the subscript on the m refers to the direction of axial stress

$$m_x := \int_{-\frac{t}{2}}^{\frac{t}{2}} \sigma_x(z) \cdot z \, dz$$

$$m_y := \int_{-\frac{t}{2}}^{\frac{t}{2}} \sigma_y \cdot z \, dz$$

$$m_x \rightarrow \frac{1}{12} \cdot t^3 \cdot E \cdot \frac{\frac{d}{dx} \frac{d}{dx} w(x, y) + \nu \frac{d}{dy} \frac{d}{dy} w(x, y)}{-1 + \nu^2}$$

$$m_y \rightarrow \frac{1}{12} \cdot t^3 \cdot E \cdot \frac{\frac{d}{dy} \frac{d}{dy} w(x, y) + \nu \frac{d}{dx} \frac{d}{dx} w(x, y)}{-1 + \nu^2}$$

with

$$D := \frac{E \cdot t^3}{12 \cdot (1 - \nu^2)}$$

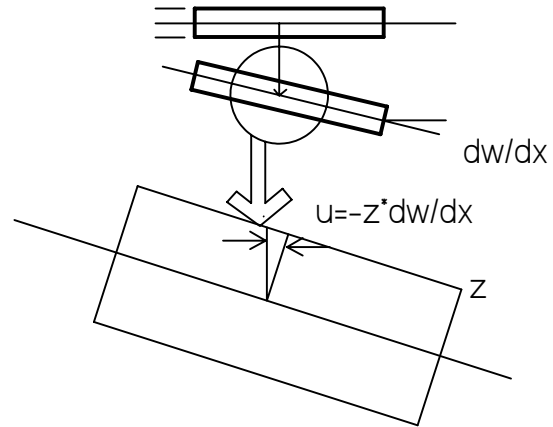
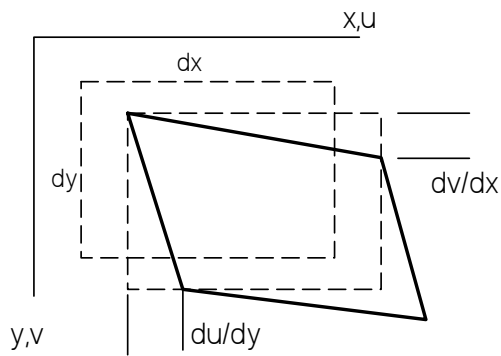
$$m_x := -D \cdot \left[\frac{d^2}{dx^2} w(x, y) + \nu \left(\frac{d^2}{dy^2} w(x, y) \right) \right]$$

$$m_y := -D \cdot \left(\frac{d^2}{dy^2} w(x, y) + \nu \frac{d^2}{dx^2} w(x, y) \right)$$

in the general plate case there may exist a twisting moment m_{xy}

we need to derive a similar expression for this moment

from basic shear strain relationships: $\gamma := \frac{d}{dx} v + \frac{d}{dy} u$; $\tau := G \cdot \gamma$ and $\tau := G \cdot \frac{d}{dx} v + \frac{d}{dy} u$



from the geometry of the slope of w in each direction (x and y): $u := -z \cdot \frac{d}{dx} w(x,y)$ and $v := -z \cdot \frac{d}{dy} w(x,y)$

$\Rightarrow \tau := -2 \cdot G \cdot z \cdot \frac{d}{dx} \frac{d}{dy} w(x,y)$ and the twisting moment per unit length is determined by:

$$m_{xy} := - \int_{-\frac{t}{2}}^{\frac{t}{2}} \tau \cdot z \, dz \quad m_{xy} \rightarrow \frac{1}{6} \cdot t^3 \cdot G \cdot \frac{d}{dx} \frac{d}{dy} w(x,y)$$

minus sign comes from sense of m_{xy} + and $\tau_{xy}(+) \cdot z(+)$ see figure Hughes 9.3 below

using: $G := \frac{E}{2 \cdot (1 + \nu)}$ $\Rightarrow m_{xy} := \frac{d}{dx} \left(\frac{d}{dy} w(x,y) \right) \cdot \frac{E t^3}{12 \cdot (1 + \nu)}$

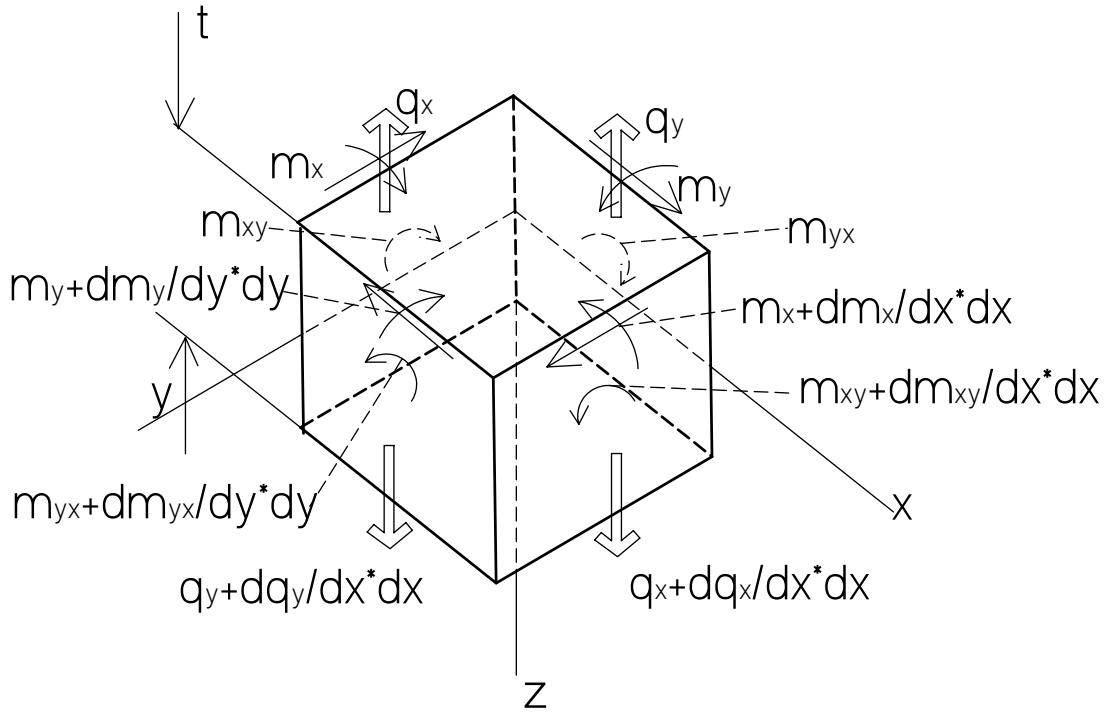
multiply and divide by $(1 - \nu)$ $\Rightarrow m_{xy} := \frac{d}{dx} \left(\frac{d}{dy} w(x,y) \right) \cdot \frac{E t^3}{12 \cdot (1 + \nu)} \cdot \frac{(1 - \nu)}{(1 - \nu)}$

$$m_{xy} := \frac{d}{dx} \left(\frac{d}{dy} w(x,y) \right) \cdot \frac{E t^3}{12 \cdot (1 - \nu^2)} \cdot (1 - \nu) \Rightarrow m_{xy} := D \cdot (1 - \nu) \left[\frac{d}{dx} \left(\frac{d}{dy} w(x,y) \right) \right]$$

due to the complimentary nature of shear stress $\tau_{xy} := \tau_{yx}$
and due to the sign convention for + moment $m_{xy} := -m_{yx}$

$$\Rightarrow m_{yx} := -D \cdot (1 - \nu) \left[\frac{d}{dx} \left(\frac{d}{dy} w(x,y) \right) \right]$$

now for equilibrium of the dx, dy, dz segment:
 see Hughes figure 9.3 forces and moments in a plate (below)



Hughes figure 9.3 forces and moments in a plate

q in this context is the shear force per unit length on each face:
 $q_x \cdot dy$ is the force on the x face etc...

equilibrium of vertical forces \Rightarrow p is the lateral (+ z direction) load per unit area - not shown on figure

$$\left[q_x + \left(\frac{d}{dx} q_x \right) \cdot dx \right] \cdot dy - q_x \cdot dy + \left[q_y + \left(\frac{d}{dy} q_y \right) \cdot dy \right] \cdot dx - q_y \cdot dx + p \cdot dx \cdot dy = 0 \quad \Rightarrow \quad \frac{d}{dx} q_x + \frac{d}{dy} q_y + p = 0$$

moments wrt the x axis \Rightarrow

$$\left[m_{xy} + \left(\frac{d}{dx} m_{xy} \right) \cdot dx \right] \cdot dy - m_{xy} \cdot dy + m_y \cdot dx - \left[m_y + \left(\frac{d}{dy} m_y \right) \cdot dy \right] \cdot dx + q_y \cdot dx \cdot dy = 0 \quad \Rightarrow \quad \frac{d}{dx} m_{xy} - \frac{d}{dy} m_y + q_y = 0$$

and similarly: taking moments wrt the y -axis \Rightarrow

$$\frac{d}{dy} m_{yx} + \frac{d}{dx} m_x - q_x = 0$$

or ... since $m_{xy} := -m_{yx}$ $-\left(\frac{d}{dy} m_{xy} \right) + \frac{d}{dx} m_x - q_x = 0$

substituting the moment relations

$$q_x := -\left(\frac{d}{dy}m_{xy}\right) + \frac{d}{dx}m_x \quad \text{and} \quad q_y := -\left(\frac{d}{dx}m_{xy}\right) + \frac{d}{dy}m_y$$

into the shear equation $\Rightarrow \frac{d}{dx}\left[-\left(\frac{d}{dy}m_{xy}\right) + \frac{d}{dx}m_x\right] + \frac{d}{dy}\left[-\left(\frac{d}{dx}m_{xy}\right) + \frac{d}{dy}m_y\right] + p = 0$

$$\frac{d^2}{dx^2}m_x - 2 \cdot \frac{d}{dx} \frac{d}{dy}m_{yx} + \frac{d^2}{dy^2}m_y + p = 0$$

substituting the relations for m_x , m_y and m_{xy} above \Rightarrow

$$\frac{d^2}{dx^2}\left[-D \cdot \left[\frac{d^2}{dx^2}w(x,y) + \nu \left(\frac{d^2}{dy^2}w(x,y)\right)\right]\right] - 2 \cdot \frac{d}{dx}\left[\frac{d}{dy}D \cdot (1-\nu) \left[\frac{d}{dx}\left(\frac{d}{dy}w(x,y)\right)\right]\right] + \frac{d^2}{dy^2}\left[-D \cdot \left(\frac{d^2}{dy^2}w(x,y) + \nu \cdot \frac{d^2}{dx^2}w(x,y)\right)\right] = -p$$

the terms with ν cancel and the result is: $\frac{d^4}{dx^4}w(x,y) + 2 \cdot \frac{d^2}{dx^2} \frac{d^2}{dy^2}w(x,y) + \frac{d^4}{dy^4}w(x,y) = \frac{p}{D}$

