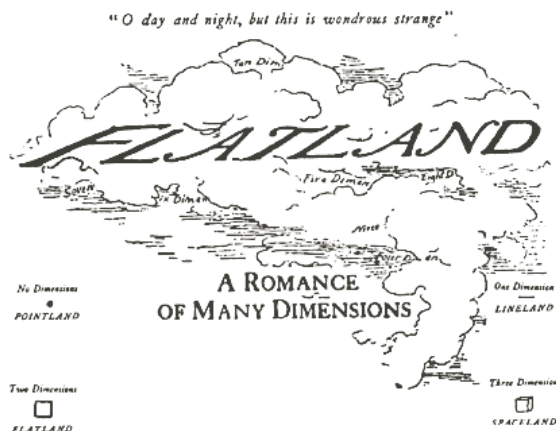


Exam #4

Problem 1 (35 points) Flatland



FLATLAND, Edwin A. Abbot, 1884

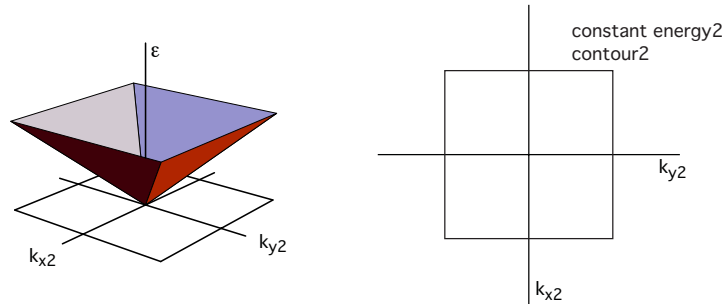
Consider world, perhaps Abbot's Flatland, where electromagnetic waves can only propagate in two dimensions, call them x and y . The electric field \vec{E} must also be in the plane, but the magnetic field \vec{B} is perpendicular to both the plane and the wavevector \vec{k} . The normal modes of the radiation field in a square box of side L with conducting walls are given by

$$\vec{E}_{k_x, k_y} = |E| \hat{1} \sin(k_x x) \sin(k_y y) \sin(\omega t + \phi)$$

where $\hat{1}$ is a unit vector in the direction of \vec{E} , k_x and k_y are determined by the need for \vec{E} to go to zero at the walls, and $\omega = c|\vec{k}|$.

- a) What are the allowed values of \vec{k} in the box?
- b) Find $D(\omega)$, the density of normal modes at frequency ω .
- c) Find an expression for $u(\omega, T)$, the temperature dependent energy density (per unit area, per unit frequency interval) of thermal radiation in this world. Do not include contributions from the zero point energy in the field.
- d) How is the Stefan-Boltzmann law changed in this world?

Problem 2 (35 points) Two-Dimensional Metal



We have studied electrons moving in a box in which the potential energy was zero. Alternatively one could consider electrons moving in a box containing a periodic potential – a simple model for the conduction electrons in a metal with a crystalline lattice. Under these conditions the single particle states can still be indexed by a wavevector \vec{k} ; however, the energy of each state $\epsilon(\vec{k})$ need not be quadratic in \vec{k} nor even isotropic in space.

The figure at the left above shows an approximation to the dispersion relation, $\epsilon(\vec{k})$, in a particular two-dimensional metal*. The energy has the form of an inverted square pyramid. It has four fold rotational symmetry. Along the k_x direction the energy is given by $\epsilon(k_x) = \gamma k_x$. The figure on the right shows a contour of constant energy on the k_x, k_y plane.

- If one imposes periodic boundary conditions on the electron wavefunctions in a square sample of side L , what are the allowed values of the wavevector \vec{k} ?
- Find $D(\vec{k})$, the density of allowed wavevectors as a function of \vec{k} .
- Find $D(\epsilon)$, the density of single particle states for the electrons as a function of their energy ϵ . Make a carefully labeled sketch of your result.
- The metal contains N conduction electrons. Find the Fermi energy ϵ_F , the energy of the last single particle state occupied at $T = 0$.
- Find the total energy of the electrons at $T = 0$ in terms of N and ϵ_F .
- Without doing any calculations, indicate how the electronic heat capacity depends on the temperature for temperatures $T \ll \epsilon_F/k_B$.
- What is the surface tension \mathcal{S} (the negative of the spreading pressure) of the electron gas at $T = 0$?

*Two dimensional planes of conduction electrons are not a fiction. They play an important role in semiconductor electronics and in high temperature superconductivity.

Problem 3 (30 points) Paramagnetic Ions

—————	①	$2\varepsilon = \mu_0 H^2$	$\mu_z = -\mu_0$
—————	②	$\varepsilon = 0$	$\mu_z = 0$
—————	①	$2\varepsilon = -\mu_0 H^2$	$\mu_z = \mu_0$

Certain impurity ions in a crystalline lattice interact with the neighboring atoms to create 4 states, 2 of which remain degenerate when a magnetic field H is applied along the z direction. The three resulting energy levels are shown above, along with their degeneracies, energies and magnetic moments.

- a) Find the partition function for a single ion, $Z_1(T, H)$. You may wish to simplify the resulting expression using hyperbolic functions; see the information sheet for the properties of the hyperbolic functions.
- b) Find the total energy $E(T, H) \equiv N \langle \varepsilon \rangle$ of N non-interacting ions in thermal equilibrium at temperature T .
- c) Find the total magnetic moment (in the z direction) due to the N ions, $M(T, H)$.

You can check your answers to b) and c) by determining if they have the expected asymptotic behavior at low and high temperature.

Work in simple systems

Hydrostatic system	$-PdV$
Surface film	$\mathcal{S}dA$
Linear system	$\mathcal{F}dL$
Dielectric material	$\mathcal{E}d\mathcal{P}$
Magnetic material	HdM

Thermodynamic Potentials when work done on the system is $dW = Xdx$

Energy	E	$dE = TdS + Xdx$
Helmholtz free energy	$F = E - TS$	$dF = -SdT + Xdx$
Gibbs free energy	$G = E - TS - Xx$	$dG = -SdT - xdx$
Enthalpy	$H = E - Xx$	$dH = TdS - xdx$

Results from hyperbolic trigonometry

$$\begin{aligned} \sinh(u) &= (e^u - e^{-u})/2 & \cosh(u) &= (e^u + e^{-u})/2 \\ \tanh(u) &= \sinh(u)/\cosh(u) & \coth(u) &= 1/\tanh(u) \\ \frac{d}{dx}(\sinh u) &= (\cosh u) \frac{du}{dx} & \frac{d}{dx}(\cosh u) &= (\sinh u) \frac{du}{dx} \end{aligned}$$

Limiting behavior of	as $u \rightarrow 0$	as $u \rightarrow \infty$
$\sinh(u)$	u	$e^u/2$
$\cosh(u)$	$1 + u^2/2$	$e^u/2$
$\tanh(u)$	u	1
$\coth(u)$	$1/u + \frac{1}{3}u$	1

Statistical Mechanics of a Quantum Harmonic Oscillator

$$\begin{aligned} \epsilon(n) &= (n + \frac{1}{2})\hbar\omega & n &= 0, 1, 2, \dots \\ p(n) &= e^{-(n+\frac{1}{2})\hbar\omega/kT} / Z(T) \\ Z(T) &= e^{-\frac{1}{2}\hbar\omega/kT} (1 - e^{-\hbar\omega/kT})^{-1} \\ \langle \epsilon(n) \rangle &= \frac{1}{2}\hbar\omega + \hbar\omega(e^{\hbar\omega/kT} - 1)^{-1} \end{aligned}$$

Radiation laws

Kirchoff's law: $e(\omega, T)/\alpha(\omega, T) = \frac{1}{4}cu(\omega, T)$ for all materials where $e(\omega, T)$ is the emissive power and $\alpha(\omega, T)$ the absorptivity of the material and $u(\omega, T)$ is the universal blackbody energy density function.

Stefan-Boltzmann law: $e(T) = \sigma T^4$ for a blackbody where $e(T)$ is the emissive power integrated over all frequencies. ($\sigma = 56.9 \times 10^{-9}$ watt-m⁻²K⁻⁴)

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8.044 Statistical Physics I
Spring 2013

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