

**PROFESSOR:** Now let me conclude for a few minutes by introducing the idea of how we're going to perturb things. So how are we going to set up our perturbation theory?

So for a perturbation theory we will do the following. We will take our system and introduce again-- so setting up the perturbation expansion. So we want to give you just an idea of what we're going to do. So  $H$  of  $t$  is going to be  $H$  of  $0$  plus  $\delta$  of  $H$  of  $t$ . And now your  $\lambda$  is going to be treated as small.

So again, we're going to use a power series expansion, and everything is going to depend on  $\lambda$ . And we're going to think coefficients in that way. We're going to work with the state  $\psi$  tilde. You remember, if you know  $\psi$  tilde, you put an  $e$  to the minus  $i\hbar 0t$  over  $\hbar$  and you can get  $H$ . So we'll set the perturbation for  $\psi$  tilde.

And it will have a zeroth part that is time dependent plus a first part that is time dependent. Every term is going to be time dependent in the perturbation. And what was our Schroedinger equation? Our Schroedinger equation was  $i\hbar \frac{d}{dt}$  of  $\psi$  tilde of  $t$  was  $\delta H$ , like that, times  $\psi$  of  $t$ .

But we need to change things a little bit. We replace  $\delta H$  by  $\lambda \delta H$ . So now it's going to have a  $\lambda$  here. So I now need to plug in this thing, which is not going to be too difficult. It's easier than what we did in time independent perturbation theory.

$\frac{d}{dt}$  of  $\psi_0$ -- let me not put the time dependence in the case  $\psi_1 \lambda^2 \psi_2$ . And this is equal to  $\lambda \delta H \psi$  tilde again. So it's  $\psi_0$  plus  $\psi_1$  plus those terms.

So for here we'll just read the first few terms. They're pretty easy. There's not much of a  $\lambda$  thing there. So what do we get? Terms without  $\lambda$ .  $i\hbar \frac{d}{dt}$  of  $\psi$  tilde  $_0$  of  $t$  equals  $0$ . This is  $\lambda$  to the  $0$ . That's the only term without the  $\lambda$ , the one that arises here.

Terms without  $\lambda$  already. Well, there's one term here, which is Schroedinger-like.  $\frac{d}{dt}$  of  $\psi_1$  is, in fact, equal to order of  $\lambda$ . You have  $\delta H$  tilde  $\psi_0$  of  $t$ .

And the next one,  $i\hbar \frac{d}{dt}$  of  $\psi_2$  of  $t$ -- that's  $\lambda^2$ -- comes from the derivative acting here. We have to look for  $\lambda^2$  here. And I forgot this  $\lambda$ . And therefore, this time you get  $\delta H \psi_1$  of  $t$ .

In general, for  $\lambda_n$ , you will get  $i\hbar \frac{d}{dt} \psi_n$ -- I'll actually put  $n+1$ ,  $n+1$  here-- is given by  $\Delta H$  acting on the previous one. So this will be simple. Once you know  $\psi_0$ , which is a constant, you put it here. This will be easily solved as an integral.

Once you have  $\psi_1$ , you put it here. You easily solve  $\psi_2$  and start solving one after another. The fun thing is that you can write these equations explicitly. And just even the first order result, and sometimes the second, give you all the physics you want, which we will explore in the next few lectures.